

Large Optical Nonlinearity of Dielectric Nanocavity-Assisted Mie Resonances Strongly Coupled to an Epsilon-near-Zero Mode

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Cite This: *Nano Lett.* 2022, 22, 702–709



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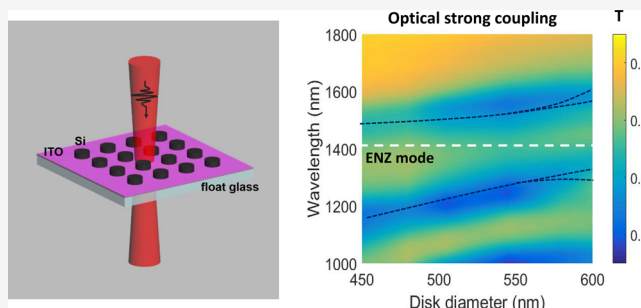


Article Recommendations



Supporting Information

ABSTRACT: Strong coupling provides a powerful way to modify the nonlinear optical properties of materials. The coupling strength of the state-of-the-art strongly coupled systems is restricted by a weak-field confinement of the cavity, which limits the enhancement of the optical nonlinearity. Here, we investigate a strong coupling between Mie resonant modes of high-index dielectric nanocavities and an epsilon-near-zero mode of an ultrathin indium tin oxide film and obtain an anticrossing splitting of 220 meV. Static nonlinear optical measurements reveal a large enhancement in the intensity-independent effective optical nonlinear coefficients, reaching more than 3 orders of magnitude at the coupled resonance. In addition, we observe a transient response of ~ 300 fs for the coupled system. The ultrafast and large optical nonlinear



coefficients presented here offer a new route towards strong coupling-assisted high-speed photonics.

KEYWORDS: strong coupling, dielectric nanocavity, Mie resonances, epsilon-near-zero mode, third-order optical nonlinearity

INTRODUCTION

Strong coupling provides a new platform for manipulating the nonlinear optical properties of materials. In such systems, an optical mode is coupled resonantly to a given optical transition (optical mode, exciton or vibration); they both exchange energy faster than each dissipative and incoherent process.^{1–4} Strong coupling can efficiently enhance many nonlinear optical phenomena of the coupled material, such as second-harmonic generation,^{5–7} third-harmonic generation,^{8,9} and four-wave mixing¹⁰ as well as ultrafast nonlinearities.^{11,12} Apart from these specific nonlinear properties, the intrinsic third-order nonlinear response, which can be described by a nonlinear refractive index n_2 and a nonlinear absorption coefficient β , has also been studied very recently in the case of strong coupling,¹³ where strongly coupled J-aggregated molecules yielded two-orders-of-magnitude enhancement factors on both n_2 and β . These strong-coupling-enhanced nonlinear coefficients facilitate the high efficiency of various third-order nonlinear optical phenomena. However, the relative weak-field confinement of the Fabry–Perot cavity used in previous work restricts the coupling strength, thereby limiting the nonlinear enhancement. One possible option to enhance the local field is to use a plasmonic cavity, considering that a plasmonic structure can confine light into a sub-wavelength volume and thereby significantly boost the coupling strength.² In a remarkable study, Boyd and coauthors showed that, indeed, very large third-order nonlinearities can be obtained when the plasmonic antennas are strongly coupled to an ultrathin indium tin oxide

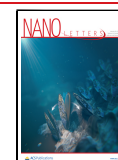
(ITO) film.¹⁴ Plasmonic nanostructures mainly enhance the electric near field while affecting only weakly the magnetic near field.¹⁵ This diminishes the technological potential of such systems in the context of nonlinear optics, including light frequency conversion and optical modulation.¹⁶

High-index dielectric nanostructures appear promising as alternative candidates to provide optical modes required for strong coupling. Indeed, dielectric resonant nanocavities have various advantages, such as low optical loss, high radiation efficiency, and reasonable bandwidth. Moreover, dielectric nanostructures with Mie-type resonances can sustain both electric dipolar (ED) and magnetic dipolar (MD) modes with comparable strength.¹⁵ In this case, the magnetic dipole stems from the circular displacement currents within each dielectric resonator induced by the incoming light. This makes it possible to enhance both the electric and magnetic near-fields at the nanoscale.^{15,17} Considering their large breakdown threshold to laser power, dielectric nanostructures are expected to be an outstanding platform for various nonlinear optical phenomena, including harmonic generations^{18–23} and ultrafast optical switching.^{24–26}

Received: October 7, 2021

Revised: December 22, 2021

Published: January 7, 2022



Taking advantage of the remarkable nonlinear optical properties of the high-index dielectric nanostructures and epsilon-near-zero (ENZ) media,^{27–29} here, we first present evidence of strong coupling between spectrally overlapping ED and MD modes of a dielectric nanocavity and an ENZ mode from an ultrathin ITO layer, reaching an anticrossing splitting as large as 220 meV. Then, by making use of the Z-scan technique,³⁰ we observe large values of effective optical nonlinear coefficients of the system at the coupled resonance, which is more than 3 orders of magnitude larger than that of bare ITO and flat silicon films. These large optical nonlinearities result from the modified properties of the coupled system, which are related to the strong local field in these structures. In addition, the coupled system exhibits a femtosecond response time of ~ 300 fs that stems from the thermalization of pump-induced hot electrons in both constituents. The large, ultrafast nonlinear response of this system is a very promising path for nonlinear photonics.

RESULTS AND DISCUSSION

Linear Optical Measurements. The coupled system in our study consists of an amorphous silicon disk array, which was fabricated on top of a 33 nm ultrathin ITO layer, as shown in Figure 1a. The silicon disks, with an average height of 200

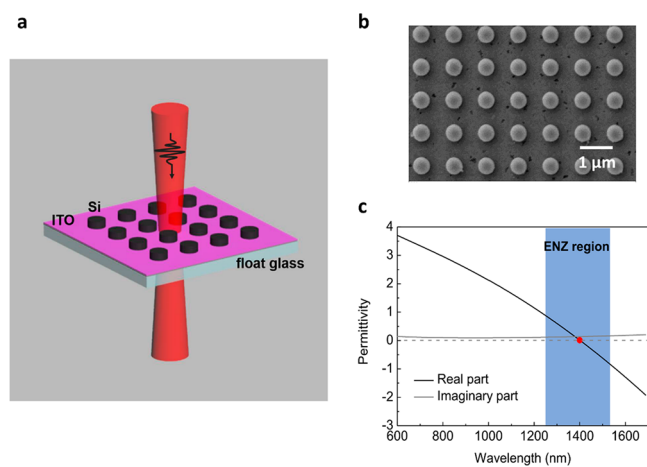


Figure 1. Schematic of the coupled system and permittivity of the ITO. (a) Illustration of the coupled system between silicon disks and ITO film. (b) SEM image of the silicon disks on top of an ITO layer. (c) Measured real (black curve) and imaginary (gray curve) parts of the permittivity of the ITO film used in this work. The ENZ region and ENZ wavelength are indicated by a blue shadow and a red dot, respectively.

nm, were fabricated on a sputtered silicon film with a standard electron-beam lithography (EBL) technique and a reactive ion etching (RIE) technique (Figure 1b). The dimensions of the unit cell were $1 \mu\text{m} \times 1 \mu\text{m}$ in order to ensure a strong far-field signal in transmittance measurements. The details of the sample's preparation are presented in the Methods section. To systematically explore the strong coupling between the ED/MD resonant modes and the ENZ mode, the diameter of the silicon disk was intentionally designed to support localized ED/MD modes at the wavelengths corresponding to the ENZ region of the ITO (Figure 1c). The linear permittivity of the ITO layer was measured by a spectroscopic ellipsometry (M-2000, J.A. Woollam Co. Inc.), and the corresponding zero

permittivity (real part) occurs at 1400 nm, as presented in Figure 1c.

In order to explore the linear optical properties of the strongly coupled ENZ mode with the Mie resonances, it is necessary to briefly discuss the optical behavior of each constituent. The ENZ mode is a confined optical mode and can only be supported by an ultrathin layer of ENZ materials.^{31–33} Its dispersion relation can be derived by solving Maxwell's equations of a three-layer system, as follows³³

$$1 + \frac{\epsilon_0 k_{z2}}{\epsilon_2 k_{z0}} = i \tan(k_{z1} d) \left(\frac{\epsilon_1 k_{z2}}{\epsilon_2 k_{z1}} + \frac{\epsilon_0 k_{z1}}{\epsilon_1 k_{z0}} \right) \quad (1)$$

where $k_{zi}^2 = \epsilon_i \frac{\omega^2}{c^2} - k_{\parallel}^2$ and ϵ_i are the longitudinal wave-number and the relative permittivity for free space ($i = 0$), ENZ material ($i = 1$), and glass substrate ($i = 2$), respectively. k_{\parallel} is the transverse wavenumber, ω is the angular frequency, d is the film thickness of the ENZ material, and c is the speed of light in free space. When the film thickness is much thinner than the skin depth of the ENZ material, such as the 33 nm ITO used in our study, the dispersion of the ENZ mode will be close to the ENZ frequency and stay nearly constant at a small k_{\parallel} region, as shown in Figure 2a. Since the dispersion of the ENZ mode lies on the right side of the air-light line, a hybrid structure that can support a wide range of wavevector components is required to efficiently excite the ENZ mode.³⁴ In addition, the electric field intensity distribution of the ENZ mode can be calculated by using field-continuity conditions (inset of Figure 2a). Obviously, the field is concentrated inside the film, which enables a high coupling efficiency when the ENZ material interacts with adjacent nanostructures.

Regarding the high-index dielectric nanostructure, a disk shape was used here because it can bring the lowest-order electric and magnetic dipolar resonances into a spectral overlap within the ENZ region by choosing a suitable height-to-diameter aspect ratio.³⁵ Therefore, it is sufficient to change the diameter of the disk while keeping the height and unit cell dimensions fixed in our simulations (COMSOL Multiphysics). The calculated transmittance spectra of the silicon disk array on a glass substrate (no ITO layer) are presented in Figure 2b. The overlapping ED and MD modes redshift as the diameter of the silicon disk increases, and eventually spectrally split when the disk diameter is larger than ~ 570 nm. This spectral separation takes place because the electric dipole resonance exhibits a much stronger redshift with respect to the nanodisk diameter as compared to the spectral shift of the magnetic dipole resonance based on the magnetic wall boundary condition of the cylindrical dielectric cavity.³⁶ The result agrees well with the diameter-dependent transmittance spectra of the silicon disks that have been discussed in the literature.³⁵ The parameters of the silicon disks are intentionally selected here so that both the ED and MD modes can appear at the ENZ region of the ITO.

The linear optical response of the silicon disk-ITO strongly coupled systems were recorded at normal incidence (model V-670, JASCO). As shown in Figure 2c, the measured transmittance spectra of the coupled systems at various disk diameters exhibit pronounced splittings at 1400 nm as well as an anticrossing feature. This feature is the signature of the strong coupling between the ENZ mode and nanodisk resonant modes.^{2,34} When the diameter of the silicon disk increases to 580 nm, the splitting of the ED and MD

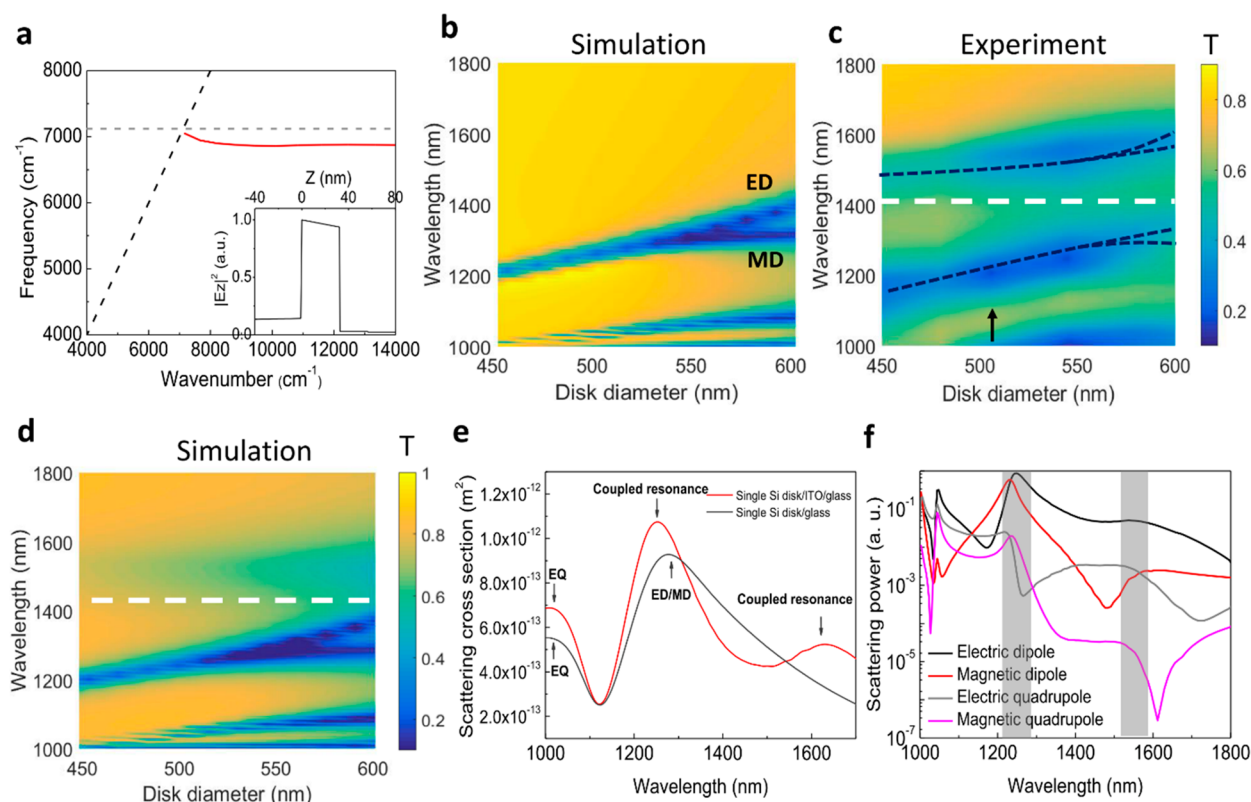


Figure 2. Linear optical responses of the coupled systems. (a) The dispersion relation of the ENZ mode (red curve) in air/33 nm ITO/glass system calculated from eq 1. The black dashed line is the air-light line $\omega = ck_{\parallel}$, and the gray dashed line denotes the ENZ frequency. (inset) Normalized electric field intensity distribution of the 33 nm ITO/substrate system. (b) Simulated transmittance spectra of silicon disks (without ITO) with respect to disk diameter. (c, d) Measured and the simulated transmittance spectra of silicon disk-ITO coupled systems as a function of disk diameter, respectively. The white dashed line represents the ENZ mode, and the black curves follow the peaks of the coupled resonances. The measured transmittance spectra at different silicon diameters are shown in Figure S1 in the Supporting Information. Here the step size in the simulations of the disk diameter is 15 nm. (e) Calculated scattering cross section of a single silicon disk (505 nm disk diameter) with (red curve) and without ITO layer (gray curve). (f) Multipolar decomposition of the 505 nm Si disk/ITO strongly coupled system with respect to wavelength. Gray areas correspond to the two strongly coupled hybrid resonances.

components at both coupled resonances occurs, which is consistent with the separation of two modes at large disk diameters in pure silicon disks as given in Figure 2b (details are shown in the Supporting Information S1). Interestingly, there are also resonances at ~ 1000 nm, that possibly come from the electric quadrupole (EQ) and collective resonance as discussed in the Supporting Information S2. Similar characteristics are also displayed in the simulated spectra of the coupled systems (Figure 2d). In contrast to the wavelength position of the resonances, the linewidths and amplitudes of the measured spectra are different from that of the simulation. These differences mainly come from the imperfect disk shape and height unavoidable through the fabrication process, because the simulated resonances are sensitive to the shape and height of the silicon disk when the disk diameter is fixed. Another possible reason is the difference between the refractive index of each material used in our simulations and actual values in the experiment. The strong coupling can also be proved by comparing the scattering cross section (SCS) spectrum of a single silicon disk with and without ITO layers (Figure 2e); here, only localized resonances exist, and there is no collective mode from grating effects. Considering a single Si disk/substrate system with a 505 nm diameter disk, the overlapping ED/MD modes appear at 1290 nm, which are close to the ENZ wavelength. When an ultrathin ITO layer is inserted between the single silicon disk and the substrate, two new

coupled resonances (one red branch and one blue branch with respect to the ENZ wavelength) are induced due to strong coupling. The coupled resonances in the SCS spectrum resemble that in the measured transmittance spectrum (blue curve in Figure S1). A multipolar decomposition analysis indicates that the electric and magnetic dipoles dominate both the coupled resonances (Figure 2f). Although the electric quadrupole and magnetic quadrupole also play a role in the coupled resonances, their contribution can be neglected in our discussion due to their small amplitudes. This suggests that both the ED and MD modes are strongly coupled to the ENZ mode.

Nonlinear Optical Measurements. The remarkable optical nonlinear response of the ENZ material and high-index dielectric nanostructures led us to examine the third-order optical nonlinearity of the strongly coupled system via Z-scan measurements. The details of our Z-scan setup were presented in the Supporting Information of our previous work.¹³ Usually, the third-order optical nonlinear process can be shown by a linear relation between the peak and valley contrast of the normalized transmittance (ΔT_{p-v}) and the peak intensity of the incident light at the focus.³⁰ In our strongly coupled system, ΔT_{p-v} increases nearly linearly with the peak intensity at a representative wavelength (Supporting Information S3). This indicates that the transmittance variation in Z-scan measurements can be used to characterize the nonlinear

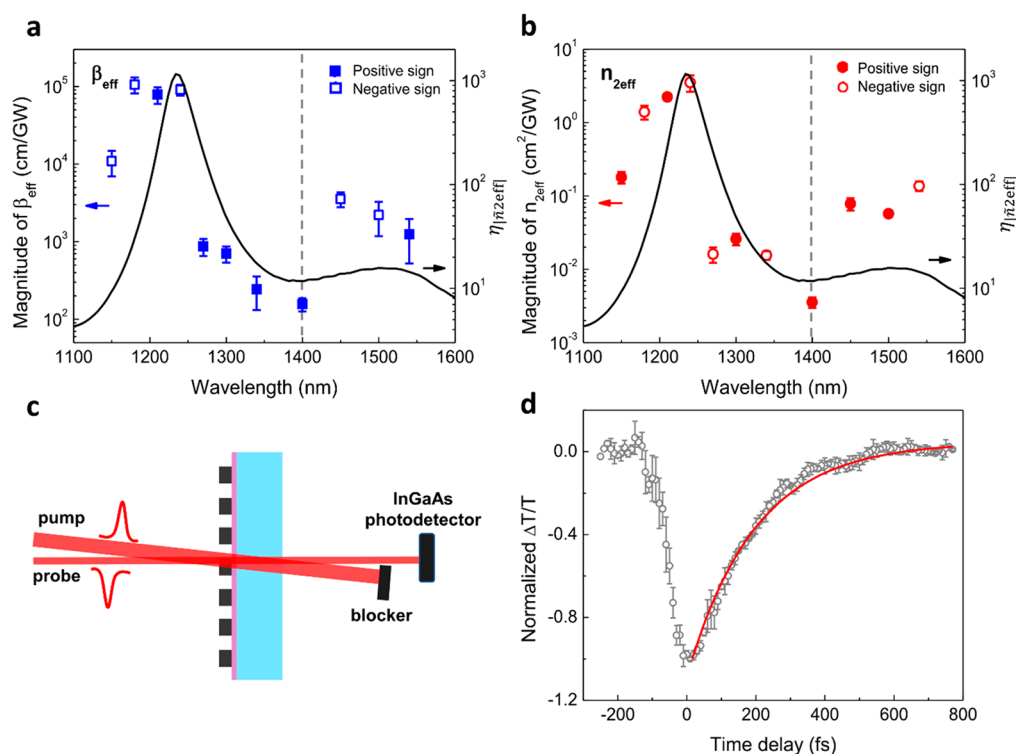


Figure 3. Nonlinear optical responses of the strongly coupled system. Magnitude of (a) β_{eff} and (b) $n_{2\text{eff}}$ for the strongly coupled system with 505 nm Si disks on an ITO film at various wavelengths. The positive and negative signs of the value are denoted as closed and open symbols, respectively. The black curves are the calculated dispersive enhancement factor of the effective complex nonlinear refractive index for the coupled system. The dashed gray lines denote the ENZ wavelength. (c) Schematic of the pump–probe experimental setup. (d) Normalized ultrafast transmittance change of the silicon disk/ITO coupled system under an 1240 nm light pulse illumination. The red curve is the fitting curve for the decay process. The decay time is 200 fs. The pump fluence applied here is 110 $\mu\text{J}/\text{cm}^2$.

refractive index and nonlinear absorption coefficient because of its third-order nonlinear origin. Here, the effective nonlinear coefficients β_{eff} and $n_{2\text{eff}}$ are used to describe the third-order optical nonlinearity in a silicon disk/ITO strongly coupled system. This way, we can estimate β_{eff} and $n_{2\text{eff}}$ from the traces of the open- and closed-aperture Z-scan, respectively (Supporting Information S4). Since both β_{eff} and $n_{2\text{eff}}$ are strongly dependent on the wavelength (Figure 3a,b), the peak intensity required for measuring a transmittance change at wavelengths far from resonances will easily cause a saturation for the wavelengths close to the resonances. To avoid this issue, lower peak intensities were used when measuring optical nonlinearity near the coupled resonances. The corresponding parameters of the laser used for the measurements at various wavelengths are given in Supporting Information S5.

In order to explore the nonlinear response under optical strong coupling conditions, we choose a silicon disk/ITO coupled sample with the largest anticrossing splitting (220 meV), which corresponds to a disk diameter of 505 nm (blue curve in Figure S1). As shown with the log plots in Figure 3a,b, both β_{eff} and $n_{2\text{eff}}$ display very large absolute values at the coupled resonances. By contrast, they exhibit relatively small values near the ENZ wavelength of the ITO. The trends of the dispersive nonlinear coefficients follow the linear transmittance spectrum of the coupled system, and the spectral positions of the maximum nonlinear coefficients match that of the coupled hybrid resonances (blue curve in Figure S1). More precisely, near the wavelength of one of the coupled resonances (1240 nm), β_{eff} and $n_{2\text{eff}}$ reach the maximum absolute values of 9.41×10^4 cm²/GW and 3.51 cm²/GW, which are more than 3 orders

of magnitude larger than both the β of a 200 nm thick silicon film (17.8 cm²/GW) and the maximal n_2 of the bare 33 nm ITO at its ENZ wavelength (1.71×10^{-3} cm²/GW). At the wavelength of another coupled resonance (~ 1550 nm), the effective nonlinearity coefficients are also enhanced by ~ 2 orders of magnitude compared with those of the bare ITO and flat silicon films. Because the optical nonlinearity of the ITO layer at the ENZ wavelength is very small compared with that of the coupled system, it is important to explore the role of the ED/MD modes of the silicon disk on this optical nonlinearity enhancement. We investigated the open- and closed-aperture Z-scan measurements of a reference sample where a silicon-disk array was directly fabricated on a float glass substrate (no ITO layer). The Si disk/substrate sample has overlapping ED/MD resonances (~ 1240 nm, Figure S6a) that are quite close to one of the coupled resonances in the Si disk/ITO/substrate system used in our optical nonlinearity measurements (blue curve in Figure S1). The values of β and n_2 of the Si-disk/substrate at the ED/MD resonant wavelength are measured to be 9.46×10^3 cm²/GW and 0.42 cm²/GW, respectively, which are still much smaller than that of the Si disk/ITO/substrate system at the coupled resonance (see Figure S6). This confirms that the strong-coupling-assisted new mode plays a major role on the large nonlinearity reported here. The Z-scan results of the bare ITO and the flat silicon films as well as the Si-disk array/substrate (without ITO) are discussed in the Supporting Information S6. In addition, the optical nonlinear coefficients at the coupled resonances measured here are also more than 2 orders of magnitude larger than that of reported plasmonic metamaterial,³⁷ and more than 1 order of

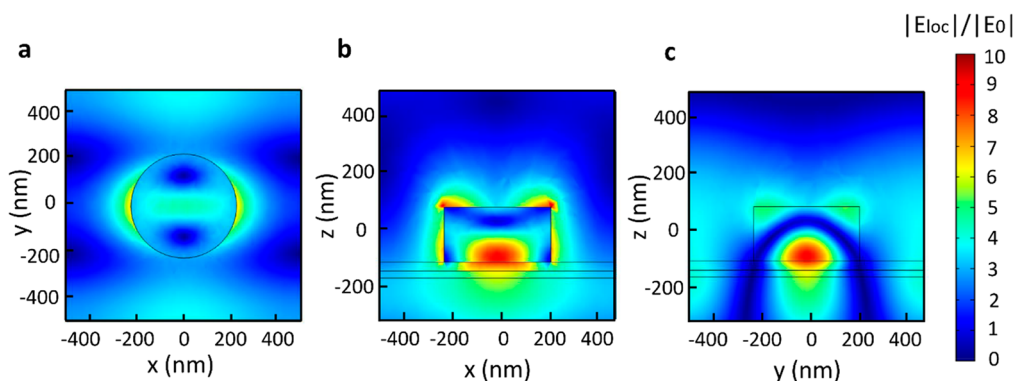


Figure 4. Electric-field distribution of the Si disk/ITO system simulated at the coupled resonant wavelength of 1240 nm. (a–c) Electric field distributions at different cross-sectional views of the silicon disk/ITO coupled system when the disk diameter is fixed to 505 nm. Here the polarization direction of the normal incident light is parallel to the x -axis.

magnitude larger than that of the nonlinear ITO in the ENZ region³⁸ and than that of a recently reported J-aggregated molecule in the strong coupling regime.¹³ The signs of the measured values of β_{eff} and $n_{2\text{eff}}$ with respect to the optical wavelength presented in Figure 3a,b are important, the negative and positive values of β_{eff} correspond, respectively, to the saturable absorption and reverse saturable absorption, while the positive and negative values of $n_{2\text{eff}}$ denote the positive and negative “Kerr lens”, respectively.

In contrast to the static optical nonlinear measurements recorded by the Z-scan technique, the time-resolved dynamics of the strongly coupled system was also exploited by modulating the amplitude of the transmitted probe light via a control beam. The schematic of the measurements is presented in Figure 3c, and the details of the pump–probe setup can be found in the Methods section. As shown in Figure 3d, with the illumination of the pump pulse, the transmittance of the probe light first reduces drastically within 100 fs near zero delay, which is restricted by the pulse width of the laser. The transmittance then recovers back to the original state within an ~ 200 fs exponential decay. The negative value of the normalized transmittance change corresponds to a light-induced absorption process for the coupled system, and the femtosecond Kerr-type nonlinearity results from the laser-induced thermalization of the hot electrons inside both the silicon disk and ITO film.^{24,38,39} The 300 fs switching time achieved here, which corresponds to a calculated bandwidth of 1.46 THz for Gaussian pulses with a time-bandwidth product of 0.44,⁴⁰ is even shorter than the optical modulation in the plasmonic antenna/ITO coupled system²⁸ that we recently reported, making our silicon disk/ITO coupled system ideal for ultrafast optical processing and switching applications.

As we stressed above, we attribute the large nonlinear coefficients at the coupled resonances to the strong electromagnetic field enhancement of the silicon disk/ITO strongly coupled system. In order to correlate the optical nonlinear coefficients and the local field intensity (I_{loc}), we analyze the electric field distribution at the coupled resonances by describing the electric field inside both the dielectric nanodisk and the ITO layer as a linear superposition of electric and magnetic components.^{41,42} Here we define an enhancement factor for the local electric field intensity as $\eta_I = I_{\text{loc}}/I_0$, where I_0 is the intensity of the incident light in vacuum. As we already explained in our former work,^{13,43} the enhancement of the effective third-order susceptibility ($\chi_{\text{eff}}^{(3)}$) and η_I can be directly related by a mean field approximation model within different

nanocavity systems. This approach leads to an enhancement factor of $\chi_{\text{eff}}^{(3)}$ that follows $\eta_{\chi_{\text{eff}}^{(3)}} = \eta_I^2$. Considering the local electric field distribution (E_{loc}) at one of the coupled resonances (1240 nm) as presented in Figure 4, η_I is estimated to be ~ 34 by averaging an unit volume of both the silicon disk and ITO layer based on the simulated electric-field profile. This gives rise to a three-orders-of-magnitude enhancement on $\chi_{\text{eff}}^{(3)}$ at the resonant condition. A comparison with the electric-field distribution at an off-resonant wavelength (1120 nm) can also be found in Figure S7 in the Supporting Information. With regard to the optical Kerr effect induced by a single light beam, the intensity-independent effective complex nonlinear refractive index $\tilde{n}_{2\text{eff}}$ follows^{44,45}

$$\tilde{n}_{2\text{eff}} = \frac{3}{4n_{0\text{eff}}n'_{0\text{eff}}\epsilon_0 c} \chi_{\text{eff}}^{(3)} \quad (2)$$

where $\tilde{n}_{2\text{eff}} = n_{2\text{eff}} + in''_{2\text{eff}} = n_{2\text{eff}} + i\frac{c}{2\omega}\beta_{\text{eff}}n_{0\text{eff}}$ and $n'_{0\text{eff}}$ are the complex and the real part of the effective linear refractive index, respectively. $n''_{2\text{eff}}$ is the imaginary part of the effective nonlinear refractive index. The relation clearly suggests that the effective nonlinear refractive index can be boosted in the same proportion as the enhancement of $\chi_{\text{eff}}^{(3)}$. Therefore, by simulating the electric-field intensity enhancement at various wavelengths (see Figure S8), the dispersive $\eta_{\chi_{\text{eff}}^{(3)}}$ and the wavelength-dependent enhancement of $\tilde{n}_{2\text{eff}}$ ($\eta_{\tilde{n}_{2\text{eff}}}$) can be estimated. Both are plotted as black curves in Figure 3a,b. As expected, the simulated enhancement of the effective nonlinear refractive index well follows the dispersive trend of the effective optical nonlinear coefficients of the strongly coupled system. We note that eq 2 can be applied here because, in our system, the coupled resonant wavelengths are relatively far from the ENZ wavelength (1400 nm).^{46,47} When considering an ENZ-effect dominated system, for instance, a thick ITO film, the free carrier nonlinearity should also be considered, as discussed in detail in the literature.^{48–53}

Although the large optical nonlinear coefficients of the silicon disk/ITO strongly coupled system are estimated from the data acquired under a strong laser irradiance, the enhancement is independent of the peak power of the incident light (Figure S3d); that is, it is an intrinsic optical property of the coupled system.

CONCLUSION

We have observed the signature of optical strong coupling between the ED/MD modes of the silicon nanodisk and the

ENZ mode of an ultrathin ITO layer. Through wavelength-dependent Z-scan measurements performed on this coupled system, we obtain large values on both $n_{2\text{eff}}$ ($-3.51 \text{ cm}^2/\text{GW}$) and β_{eff} ($-9.41 \times 10^4 \text{ cm/GW}$) at the coupled resonant wavelength. These large nonlinearities are ascribed to a significant enhancement of the electric field in the coupled system. In addition, the coupled system demonstrates an ultrafast modulation of its optical response of $\sim 300 \text{ fs}$.

Since the optical nonlinearity coefficients of the coupled system are strongly related to the local electric field enhancement, designing a specific dielectric nanocavity with a higher field confinement could be a direct way to further improve the values of $n_{2\text{eff}}$ and β_{eff} for the system. Because the excitons of semiconductors also show remarkable optical nonlinearities,⁵⁴ one can expect a large optical nonlinearity by coupling the excitons to the ED/MD modes of the dielectric nanocavity. The study of ultrafast, large third-order optical nonlinearity in a dielectric nanodisk/ITO coupled system here offers a new route to explore various third-order optical nonlinear phenomena via strong coupling for high-speed active photonic and optoelectronic devices.

METHODS

Sample Preparation. The sample used in our study is a silicon nanodisk array, which was fabricated on top of a commercially 33 nm ITO layer (sheet resistance $50 \Omega/\text{square}$) via standard EBL and RIE techniques. Between the ITO film and the substrate, there is an $\sim 25 \text{ nm}$ SiO_2 passivation layer in order to ensure electric insulation. Before fabrication, an $\sim 200 \text{ nm}$ thick amorphous silicon film was sputtered on top of the ITO layer (0.21 nm/s , PLASSYS MP300). For the EBL process, two layers of the positive Poly(methyl methacrylate) (PMMA) e-beam resists (system 600K/950 K of AR-P series, Allresist GmbH) were used in order to improve the accuracy of the exposed pattern. The 600-K and 950-K resists were spin-coated sequentially at 4000 and 2000 rpm, respectively, on top of the silicon film. The film thickness of the 600-K and 950-K resists are 220 and 100 nm, respectively. After baking the spin-coated resists at $180 \text{ }^\circ\text{C}$ for 90 s, the designed areas were exposed to a 30 keV electron beam (AURIGA, Zeiss). The exposed part of the resist was then developed in a conventional solution of methyl isobutyl ketone–isopropyl alcohol (MIBK-IPA) (1:3) for 30 s and then rinsed with an IPA solution for 30 s to complete the development process. Next, the Cr mask was obtained by evaporating a 30 nm chrome (Cr) film (0.05 nm/s , PLASSYS ME300) and removing the unexposed resist with acetone (lift-off process). The Cr disks can then be transferred to the silicon film by using Ar and CHF_3 plasma gases under the pressure of $2 \times 10^{-5} \text{ Torr}$ in a RIE chamber (PlasmaLab 80 Plus, Oxford Instruments). Finally, the Cr mask was removed by immersing the sample in a chromium etchant solution in $\sim 2 \text{ min}$. The silicon disks, which possess average diameters from 450 to 600 nm, were designed so that they support localized ED and MD modes at wavelengths that can strongly couple to the ENZ mode of ITO at 1400 nm. The choice of the unit cell dimensions is also crucial because the wavelength of collective resonances from the metasurface rely on the periodicity of the structure. Here we choose $1 \mu\text{m} \times 1 \mu\text{m}$ unit cell dimensions in order to minimize the influence of the collective resonances in the ENZ region of the ITO.

Time-Resolved Dynamics Measurements. The temporal response of the coupled system was performed with a degenerate optical pump–probe setup. The 1000 Hz, 1240 nm

light applied here was generated from an optical parametric amplifier (OPA, TOPAS), which is seeded by a Ti:sapphire laser system (Spitfire, Spectra-Physics). The output beam was split into two portions as the pump and probe beams. The two beams were set to be nearly parallel ($\sim 3^\circ$ intersection angle) and illuminated the sample surface normally. The transient transmitted intensity of the probe beam was acquired by a biased indium gallium arsenide (InGaAs) photodetector (DET10N2, Thorlabs, Inc.). The time delay between the pump and the probe pulses was set by routing the probe beam through a motorized stage with retroreflectors.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.nanolett.1c03876>.

Linear transmittance spectra of the strongly coupled systems; periodicity-dependent transmittance spectra of the strongly coupled system; intensity evolutions of open- and closed-aperture Z-scan traces for the strongly coupled system; estimation of optical nonlinear coefficients from Z-scan traces; parameters of the laser for Z-scan measurements; nonlinear coefficients of the reference samples; electric field intensity at various wavelengths (PDF)

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Author Contributions

K.W., H.-H.H., and T.W.E. conceived the idea; A.-Y.L. and H.-H.H. designed the nanostructures and performed the numerical simulations. K.W. fabricated the samples and conducted all the optical measurements and the corresponding data analysis. All authors contributed to analyzing, interpreting the experimental results, and commenting on the manuscript. C.G. and T.W.E. supervised the project.

Funding

This work of the Interdisciplinary Thematic Institute QMat, as part of the ITI 2021–2028 program of the University of Strasbourg, CNRS and Inserm, was supported by IdEx Unistra (ANR-10-IDEX-0002), by SFRI STRAT'US project (ANR-20-SFRI-0012), by the Labex NIE (ANR-11-LABX-0058 NIE), CSC (ANR-10-LABX-0026 CSC), the EQUIPEX UNION

(ANR-10-EQPX-52–01) and PlasHybrid (ANR-18-CE30–0014–02) projects, and by the University of Strasbourg Institute for Advanced Study (USIAS) (ANR-10-IDEX-0002–02) under the framework of the French Investments for the Future Program. The support of the International Center for Frontier Research in Chemistry (icFRC, Strasbourg), the ERC (Project No. 788482 MOLUSC) and QuantERA (Project RouTe) together with the Ministry of Science and Technology of Taiwan under Grant Nos. MOST 107-2112-M-003-013-MY3 and MOST 110-2112-M-003-020-MY3 are also acknowledged.

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENTS

The authors thank Hicham Majjad for the help of reactive ion etching fabrication, Yifan Yao for the use of the atomic force microscopy, and Yuanyuan Zhang for the help of permittivity measurements of ITO with the spectroscopic ellipsometry in East China Normal University. We are also thankful for fruitful discussions with Minghao Li and Gian Lorenzo Paravicini-Bagliani.

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